

A hp - discontinuous Galerkin method for the time-dependent Maxwell's equation : a priori error estimate

C. DAVEAU, A. ZAGHDANI
Département de Mathématiques
Université de Cergy-Pontoise
2 avenue Adolphe Chauvin
95302 Cergy-Pontoise cedex, France
 christian.daveau@math.u-cergy.fr

Abstract. A discontinuous Galerkin method for the numerical approximation for the time-dependant Maxwell's equations in "stable medium" with supraconductive boundary, is introduced and analysed. its hp -analysis is carried out and error estimates that are optimal in the meshsize h and slightly suboptimal in the approximation degree p are obtained.

Keywords: Discontinuous Galerkin method, wave equation, a priori error estimate.

AMS subject classifications. 65N30.

1. Introduction

Electromagnetism is one application of the discontinuous Galerkin method, among many other areas. In [1], the discontinuous Galerkin method with solutions that are exactly divergence-free inside each element, is developped for numerically solving the Maxwell equations. In [2], M. Grote, A. Schneebeli and D. Schötzau propose and analyse the symmetric interior penalty discontinuous Galerkin method for the spatial discretization of Maxwell's equations in second order form.

Here, we consider a nonsymmetric interior penalty discontinuous Galerkin method to approximate in space an initial boundary value problem derived from Maxwell's equations in vacuum with perfect electric conductor boundary

$$\frac{\partial^2 u}{\partial t^2} + c^2 \nabla \times (\nabla \times u) = f, \quad \nabla \cdot u = 0 \quad \text{in } \Omega \times I; \quad (1.1)$$

$$n \times u(x, t) = 0 \quad \text{on } \partial\Omega \times I, \quad u(x, 0) = u_0(x), \quad \frac{\partial u}{\partial t}(x, 0) = u_1(x) \quad \text{on } \Omega. \quad (1.2)$$

Here Ω is a convex polyhedron included in \mathbb{R}^3 , $I = [0, t^*] \subset \mathbb{R}$, u_0 and u_1 are in $H_0(\nabla \times, \Omega) \cap H(\nabla \cdot 0, \Omega)$ and f is defined on $\Omega \times I$. Physically, u is the electric field, and f is related to a current density. Moreover $\mu_0 \varepsilon_0 c^2 = 1$ where $\mu_0 \approx 4\pi 10^{-7}$ H.m⁻¹ and $\varepsilon_0 \approx (36\pi 10^9)^{-1}$ F.m⁻¹ are respectively the magnetic permeability and the electric permittivity in vacuum. If we assume that the domain Ω is "stable medium" with boundary and if u is the exact solution of (1.1) and (1.2) then u and $\nabla \times u$ belong to $H^1(\Omega)^3$. For the notations, if I is an interval, X is a function space and ϕ is a function on $\Omega \times I$ then $\|\phi\|_{L^p(I, X)}$ denotes the norm in $L^p(I)$ of the function $t \rightarrow \|\phi(\cdot, t)\|_X$. $L^p(X)$ is short for $L^p(I, X)$.

Let Π_h be a partition of Ω into tetrahedra and consider the same spaces and notations as in [5].

Faces: We define and characterise the faces of the triangulation Π_h . An interior face of Π_h is defined as the (non-empty) two-dimensional interior of $\partial K_1 \cap \partial K_2$, where K_1 and K_2 are two adjacents elements of Π_h . A boundary face of Π_h is defined as the (non-empty) two-dimensional interior of $\partial K \cap \partial\Omega$, where K is a boundary element of Π_h . We denote by F_h^I the union of all interior faces of Π_h , by F_h^D the union of all boundary faces of Π_h and by F_h the union of all faces of Π_h . Furthermore we identify F_h^D to $\partial\Omega$ since Ω is a polyhedron.

Traces: Let $H^s(\Pi_h) = \{v : v|_K \in H^s(K) \quad \forall K \in \Pi_h\}$ for $s > \frac{1}{2}$ endowed with the norm $\|v\|_{s, \Pi_h}^2 = \sum_{K \in \Pi_h} \|v\|_{s, K}^2$. Then the elementwise traces of functions in $H^s(\Pi_h)$ belongs to $TR(F_h) = \Pi_{K \in \Pi_h} L^2(\partial K)$; they are double-valued on F_h^I and single-valued on F_h^D . The space $L^2(F_h)$ can be identified with the functions in $TR(F_h)$ for which the two traces values coincide.

Traces operators: Let us introduce the following traces operators for piecewise smooth functions. First, let $w \in TR(F_h)^3$ and $e \subset F_h$. If e is an interior face in F_h^I , we denote by K_1 and K_2 the elements sharing e , by n_i the normal unit vector pointing exterior to K_i and we set $\omega_i = \omega|_{\partial K_i}, i = 1, 2$. We define the *average*, and the *tangential* and *normal jumps* of w at $x \in e$ as

$$\{\omega\} = \frac{\omega_1 + \omega_2}{2}, \quad [\omega]_T = n_1 \times \omega_1 + n_2 \times \omega_2 \quad \text{and} \quad [\omega]_N = n_1 \cdot \omega_1 + n_2 \cdot \omega_2.$$

If $e \subset F_h^D$, we set for $x \in e$

$$\{\omega\} = \omega, \quad [\omega]_T = n \times \omega \quad \text{and} \quad [\omega]_N = n \cdot \omega.$$

If $w \in H(\nabla \times, \Omega)$, then, for all $e \subset F_h^I$ the jump condition $[w]_T = 0$ holds true in $L^2(e)^3$ and if $w \in H(\nabla \cdot, \Omega)$, then, for all $e \subset F_h^I$ the jump condition $[w]_N = 0$ holds true in $L^2(e)$.

For $e \subset F_h$, we denote by $\langle \cdot, \cdot \rangle_e$ the scalar product in $L^2(e)^3$ or $L^2(e)$, furthermore if F_h^D is identified to $\partial\Omega$, we identify $\sum_{e \subset F_h^D} \langle \cdot, \cdot \rangle_e$ to $\langle \cdot, \cdot \rangle$, the scalar product in $L^2(\partial\Omega)^3$ or $L^2(\partial\Omega)$.

Finite element spaces: In order to define the average of $\nabla \times u$, we set for $s > \frac{1}{2}$, $H^s(\nabla \times, \Pi_h) := \{v : v|_K \in H^s(K)^3 \text{ and } \nabla \times (v|_K) \in H^s(K)^3, \forall K \in \Pi_h\}$. Let $p = (p_K)_{K \in \Pi_h}$ be a degree vector that assigns to each element $K \in \Pi_h$ a polynomial approximation order $p_K \geq 1$. The generic hp -finite element space of piecewise polynomials is given by

$$S^p(\Pi_h) = \{u \in L^2(\Omega) : u|_K \in S^{p_K}(K) \quad \forall K \in \Pi_h\}$$

where $S^{p_K}(K)$ is the space of real polynomials of degree at most p_K in K .

We also set $\Sigma_h := S^p(\Pi_h)^3$.

Now, fix a face $e \subset F_h$ and define the local parameters h, p by $h := \min(h_K, h_{K'})$, $p = \max(p_K, p_{K'})$ in the case of interior faces and $h := h_K$, $p = p_K$ in the case of boundary faces [4].

2. Discretization and a priori error estimate

In order to derive a weak formulation of (1)-(2), we note that formulas (1) in [5] implies that for any u with $\nabla \times u \in H(\nabla \times, \Omega)$

$$c^2(\nabla \times (\nabla \times u), v) = c^2(\nabla \times u, \nabla \times v) + a(u, v)$$

where

$$a(u, v) = c^2 \langle n \times (\nabla \times u), v \rangle - c^2 \sum_{e \subset F_h^I} \langle [v]_T, \{\nabla \times u\} \rangle_e.$$

Now, we introduce the penalty term via the form

$$J^\sigma(u, v) = \sum_{e \subset F_h^I} \langle \sigma[u]_N, [v]_N \rangle_e + \sum_{e \subset F_h} \langle \sigma[u]_T, [v]_T \rangle_e \quad u, v \in H^1(\nabla \times, \Pi_h)$$

where $\sigma := \kappa p^2/h$ is a stabilization parameter and κ is a constant supposed ≥ 1 . We also define

$$A(u, v) = c^2(\nabla \times u, \nabla \times v) + a(v, u) - a(u, v) + J(u, v),$$

$$B(u, v) = A(u, v) + J^\sigma(u, v)$$

and

$$J(u, v) = (\nabla \cdot u, \nabla \cdot v).$$

2.1. Properties of the bilinear form

Now, we introduce a norm associated with the bilinear form B and set for $u \in H^1(\nabla \times, \Pi_h)$

$$\begin{aligned} \|u\|_h^2 &= \|u\|^2 + \|c(\nabla \times u)\|^2 + \|\nabla \cdot u\|^2 + \left\| \frac{1}{\sqrt{\sigma}} \{\nabla \times u\} \right\|_{0, F_h}^2 \\ &\quad + \|\sqrt{\sigma}[u]_N\|_{0, F_h^I}^2 + \|\sqrt{\sigma}[u]_T\|_{0, F_h}^2. \end{aligned}$$

We start by studying the continuity of the bilinear forms introduced above. We have :

Proposition 2.1. $\forall v, u \in H^1(\nabla \times, \Pi_h)$ there exists a constant C independent of h and p such that

$$|A(u, v)| \leq C\|u\|_h\|v\|_h \quad \text{and} \quad |J^\sigma(u, v)| \leq C\|u\|_h\|v\|_h.$$

Proof. The proof is easily deduced from the definition of A , J^σ , $\|\cdot\|_h$ and the Cauchy-Schwarz inequality \square

In order to study the coercivity of the bilinear form B , we start by introducing the following inequality of Poincaré-Friedrichs type valid for $u \in H^1(\Pi_h)^3$.

Lemma 2.1. Let $u \in H^1(\Pi_h)^3$. Then there exists C independent of h and p such that

$$\|u\|^2 \leq C(\|c(\nabla \times u)\|^2 + \|\nabla \cdot u\|^2 + \sum_{e \subset F_h} \|\sqrt{\sigma}[u]_T\|_{0, e}^2 + \sum_{e \subset F_h^I} \|\sqrt{\sigma}[u]_N\|_{0, e}^2)$$

Proof. The proof follows immediately from Lemma 3.1 in [5] since $\kappa p^2 \geq 1$. \square

Now, the following coercivity result holds.

Proposition 2.2. There exists two constants $\alpha > 0$ and $\tilde{C} > 0$ independent of h and p such that

$$B(v, v) \geq \alpha\|v\|_h^2 + \tilde{C}J^\sigma(v, v) \quad \forall v \in \Sigma_h.$$

Proof. Let us first recall the following inverse inequality

$$\|q\|_{0, \partial K}^2 \leq C_{inv} \frac{p_K^2}{h_K} \|q\|_{0, K}^2 \quad \forall q \in S^{p_K}(K). \quad (2.1)$$

with a constant $C_{inv} > 0$, only depending on the shape regularity of the mesh. Now, Let α be an arbitrary real number and choose $v \in \Sigma_h$. Then

$$\begin{aligned} B(v, v) - \alpha\|v\|_h^2 &= (1 - \alpha)(A(v, v) + J^\sigma(v, v)) \\ &\quad - \alpha \int_{F_h} \{\nabla \times v\}^2 / \sigma ds - \alpha\|v\|^2. \end{aligned}$$

Since $\{\nabla \times v\}$ is the average of the flux at the face of two elements K_i and K_j , the corresponding integral can be split into two integrals with integrands $(\nabla \times v)_i / \sigma$ and $(\nabla \times v)_j / \sigma$, each one associated with the elements K_i or K_j respectively. Therefore, let $e \subset F_h$ and consider the integral associated with the element K . Using the inverse inequality (3), we have since $\nabla \times (\Sigma_h) \subset \Sigma_h$,

$$\int_e (\nabla \times v)^2 / \sigma ds = \frac{1}{\sigma} \|\nabla \times v\|_{0, e}^2 \leq \frac{C_{inv}}{\sigma} \frac{p_K^2}{h_K} \|\nabla \times v\|_{0, K}^2 \quad (2.2)$$

so that, selecting σ to be equal to $\kappa p^2 / h$ in (2.2), we obtain

$$- \int_e (\nabla \times v)^2 / \sigma ds \geq - \frac{C_{inv}}{\kappa} \|\nabla \times v\|_{0, K}^2.$$

In particular,

$$\begin{aligned} - \int_{F_h} \{\nabla \times v\}^2 / \sigma ds &\geq - \frac{C_{inv}}{\kappa} \sum_{K \in \Pi_h} \|\nabla \times v\|_{0, K}^2 \\ &\geq - \frac{C_{inv}}{\kappa} A(v, v). \end{aligned}$$

It then follows that

$$B(v, v) - \alpha \|v\|_h^2 \geq (1 - \alpha - \alpha \frac{C_{inv}}{\kappa}) A(v, v) + (1 - \alpha) J^\sigma(v, v) - \alpha \|v\|^2$$

and we can easily see there exists a positive constant C such that

$$\|u\|^2 \leq C(A(v, v) + J^\sigma(v, v)).$$

Then, we obtain

$$\begin{aligned} B(v, v) - \alpha \|v\|_h^2 &\geq (1 - \alpha C - \alpha - \alpha C_{inv}/\kappa) A(v, v) + (1 - \alpha - \alpha C) J^\sigma(v, v) \\ &\geq (1 - \alpha(C + 1 + C_{inv}/\kappa)) A(v, v) + (1 - \alpha(1 + C)) J^\sigma(v, v). \end{aligned}$$

Thus, if α is chosen for example $\alpha = \frac{1}{C+1+C_{inv}/\kappa}$ and $\tilde{C} = 1 - \frac{1+C}{C+1+C_{inv}/\kappa} > 0$, we immediately obtain the coercivity result. \square

Now, the following hp -approximation result to interpolate scalar function holds (see [4]).

Prop 2.1. Let $K \in \Pi_h$ and suppose that $u \in H^{t_K}(K)$, $t_K \geq 1$. Then there exists a sequence of polynomials $\pi_{p_K}^{h_K}(u) \in S^{p_K}(K)$, $p_K = 1, 2, \dots$ satisfying, $\forall 0 \leq q \leq t_K$

$$\|u - \pi_{p_K}^{h_K}(u)\|_{q,K} \leq C \frac{h_K^{\min(p_K+1, t_K)-q}}{p_K^{t_K-q}} \|u\|_{t_K, K} \quad \text{and}$$

$$\|u - \pi_{p_K}^{h_K}(u)\|_{0, \partial K} \leq C \frac{h_K^{\min(p_K+1, t_K)-\frac{1}{2}}}{p_K^{t_K-\frac{1}{2}}} \|u\|_{t_K, K}.$$

The constant C is independent of u , h_K and p_K , but depends on the shape regularity of the mesh.

In order to interpolate vector function, we define

Definition 2.1. For $u = (u_1, u_2, u_3)$ we define

$\mathbf{\Pi}_p^h : H^t(\nabla \times, \Pi_h) \rightarrow \Sigma_h$ by $\mathbf{\Pi}_p^h(u) = (\pi_p^h(u_1), \pi_p^h(u_2), \pi_p^h(u_3))$ with π_p^h is defined by $\pi_p^h(u)|_K = \pi_{p_K}^{h_K}(u|_K)$ where $\pi_{p_K}^{h_K}$ is given by the previous Proposition.

2.2. A priori error estimate

The interior penalty finite element approximation to u is to find $U : I \rightarrow \Sigma_h$ such that

$$(U_{tt}, v) + B(U, v) = (f, v) \quad \forall v \in \Sigma_h, \quad U(0) = \mathbf{\Pi}_p^h(u_0), \quad U_t(0) = \mathbf{\Pi}_p^h(u_1). \quad (2.3)$$

Upon choice of a basis for Σ_h and the data f , (2.3) determines U as the only solution to an initial value problem for a linear system of ordinary differential equations. Note that, if u is the exact solution of (1.1)-(1.2), then u satisfies the first equation in (2.3) and thus the problem is consistent.

We now analyse the proposed procedure by the method of energy estimates. In this Section, u denotes the exact solution of (1.1)-(1.2) and U the discrete solution of (2.3). C is generic constant independent of h and p which takes different values at the different places and depends on α , \tilde{C} the coercivity constants of the form B , t^* and Ω .

Let $\zeta = U - u$, then ζ satisfies

$$(\zeta_{tt}, v) + B(\zeta, v) = 0 \quad \forall v \in \Sigma_h.$$

Decompose ζ as $\mu - \nu$ where $\mu = \mathbf{\Pi}_p^h(u) - u$ and $\nu = \mathbf{\Pi}_p^h(u) - U$. Thus

$$(\nu_{tt}, v) + B(\nu, v) = (\mu_{tt}, v) + B(\mu, v) \quad \forall v \in \Sigma_h.$$

Since $\nu_t(t) \in \Sigma_h$, we can set $v = \nu_t(t)$, obtaining

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\nu_t(t)\|^2 + \frac{1}{2} \frac{d}{dt} B(\nu(t), \nu(t)) &= (\mu_{tt}(t), \nu_t(t)) + B(\mu(t), \nu_t(t)) \\ &\leq \frac{1}{2} \|\mu_{tt}(t)\|^2 + \frac{1}{2} \|\nu_t(t)\|^2 + B(\mu(t), \nu_t(t)). \end{aligned}$$

So

$$\frac{d}{dt} \|\nu_t(t)\|^2 + \frac{d}{dt} B(\nu(t), \nu(t)) \leq \|\mu_{tt}(t)\|^2 + \|\nu_t(t)\|^2 + 2B(\mu(t), \nu_t(t)).$$

Since $\nu_t(0) = \nu(0) = 0$, integration over $[0, t] \subset I$, yields

$$\|\nu_t(t)\|^2 + B(\nu(t), \nu(t)) \leq \|\mu_{tt}\|_{L^2(L^2)}^2 + \int_0^t \|\nu_t(t)\|^2 dt + 2 \int_0^t B(\mu(t), \nu_t(t)) dt.$$

The final term may be integrated by parts in time. Hence,

$$2 \int_0^t B(\mu(t), \nu_t(t)) dt \leq 2|B(\mu(t), \nu(t))| + 2 \int_0^t |B(\mu_t(t), \nu(t))| dt.$$

Therefore, we can apply the coercivity and continuity of B to get

$$\begin{aligned} &\|\nu_t(t)\|^2 + \alpha \|\nu(t)\|_h^2 + \tilde{C} J^\sigma(\nu(t), \nu(t)) \\ &\leq \|\mu_{tt}\|_{L^2(L^2)}^2 + \int_0^t \|\nu_t(t)\|^2 dt + C \|\nu(t)\|_h \|\mu(t)\|_h + 2 \int_0^t |B(\mu_t(t), \nu(t))| dt \\ &\leq \|\mu_{tt}\|_{L^2(L^2)}^2 + \int_0^t \|\nu_t(t)\|^2 dt + C \|\mu(t)\|_h^2 + \frac{\alpha}{2} \|\nu(t)\|_h^2 + C \int_0^t (\|\mu_t(t)\|_h^2 + \|\nu(t)\|_h^2) dt \\ &\leq C \left(\|\mu_{tt}\|_{L^2(L^2)}^2 + \sup_{t \in I} \|\mu(t)\|_h^2 + \int_0^{t^*} \|\mu_t(t)\|_h^2 dt \right) + \frac{\alpha}{2} \|\nu(t)\|_h^2 + C \int_0^t (\|\nu_t(t)\|^2 + \|\nu(t)\|_h^2) dt. \end{aligned}$$

In particular,

$$\begin{aligned} &\|\nu_t(t)\|^2 + \|\nu(t)\|_h^2 + J^\sigma(\nu(t), \nu(t)) \\ &\leq C \left(\|\mu_{tt}\|_{L^2(L^2)}^2 + \sup_{t \in I} \|\mu(t)\|_h^2 + \int_0^{t^*} \|\mu_t(t)\|_h^2 dt \right) + C \int_0^t (\|\nu_t(t)\|^2 + \|\nu(t)\|_h^2) dt \end{aligned}$$

As this holds for all $t \in I$, Gronwall's Lemma implies that

$$\|\nu_t(t)\|^2 + \|\nu(t)\|_h^2 + J^\sigma(\nu(t), \nu(t)) \leq C \left(\|\mu_{tt}\|_{L^2(L^2)}^2 + \sup_{t \in I} \|\mu(t)\|_h^2 + \int_0^{t^*} \|\mu_t(t)\|_h^2 dt \right).$$

Since $\zeta = \mu - \nu$ and $J^\sigma(\mu, \mu) = J^\sigma(\mu, \nu) = 0$,

$$\|\zeta_t(t)\|^2 + \|\zeta(t)\|_h^2 + J^\sigma(\zeta(t), \zeta(t)) \leq C \left(\|\mu_{tt}\|_{L^2(L^2)}^2 + \sup_{t \in I} \|\mu(t)\|_h^2 + \int_0^{t^*} \|\mu_t(t)\|_h^2 dt + \|\mu_t\|_{L^\infty(L^2)}^2 \right).$$

Thus, error bounds for the finite element approximation to the true solution reduces to the error bounds for the piecewise polynomial interpolant. Thus, we start by estimating $\|u - \mathbf{\Pi}_p^h(u)\|_h$ where $\mathbf{\Pi}_p^h$ is defined after definition 2.1. By using Proposition 2.1 and the definition of $\|\cdot\|_h$, we obtain the following estimates

$$\|u - \mathbf{\Pi}_p^h(u)\|_h^2 \leq C \sum_{K \in \Pi_h} \frac{h_K^{2\mu_K - 2}}{p_K^{2t_K - 3}} \|u\|_{t_K, K}^2 \quad \text{and} \quad \|u - \pi_{p_K}^h(u)\|_{q, K} \leq C \frac{h_K^{\mu_K - q}}{p_K^{t_K - q}} \|u\|_{t_K, K} \quad \forall 0 \leq q \leq t_K.$$

where $\mu_K = \min(p_K + 1, t_K)$. By using the previous estimates, we can get the following result

Proposition 2.3. *Let $\mu_K = \min(p_K + 1, t_K)$ and u be the exact solution of (1.1)-(1.2). Suppose that $u|_K \in C^2(I, H^{t_K}(K)^3)$, $\forall K \in \Pi_h$ with $t_K \geq 2$. Let U the discrete solution of (2.3). Then, the error $\zeta = U - u$ satisfies*

$$\begin{aligned} &\|\zeta_t(t)\|^2 + \|\zeta(t)\|_h^2 + J^\sigma(\zeta(t), \zeta(t)) \\ &\leq C \sum_{K \in \Pi_h} \frac{h_K^{2\mu_K - 2}}{p_K^{2t_K - 3}} \left(\|u_{tt}\|_{L^2(H^{t_K}(K)^3)}^2 + \|u\|_{L^\infty(H^{t_K}(K)^3)}^2 + \|u_t\|_{L^2(H^{t_K}(K)^3)}^2 + \|u_t\|_{L^\infty(H^{t_K}(K)^3)}^2 \right). \end{aligned}$$

References

- [1] B. Cockburn, F. Li, C.-W Shu, “Locally divergence-free discontinuous Galerkin method for Maxwell equations.”, *Journal of Computational Physics*, Vol. 194 , pp 588-610, 2004.
- [2] M. Grote, A. Schneebeli, D. Schötzau, “Interior penalty discontinuous Galerkin for Maxwell’s equations: Energy norm estimates.” *Journal of Computational and Applied Mathematics*, Vol. 204 , pp 375-386, 2007.
- [3] P. Houston, C. Schwab and E. Süli “Discontinuous hp -finite element methods for advection-diffusion-reaction problems. ”, *SIAM J. Numer. Anal.* 39, pp 2133–2163, 2002.
- [4] I. Perugia and D. Schötzau, “The hp -Local Discontinuous Galerkin method for the Low-Frequency Time-Harmonic Maxwell’s Equations.” *Math. Comp.*, no.243 pp 1179-1214, 2003.
- [5] A. Zaghdani and C. Daveau “Two new discrete inequalities of Poincaré-Friedrichs on discontinuous spaces for Maxwell’s equations.” *C. R. Acad. Sci. Paris. Ser.I.* 342, pp 29-32, 2006.